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# Bounds for the connective constant of the hexagonal lattice 

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#### Abstract

We give improved bounds for the connective constant of the hexagonal lattice. The lower bound is found by using Kesten's method of irreducible bridges and by determining generating functions for bridges on one-dimensional lattices. The upper bound is obtained as the largest eigenvalue of a certain transfer matrix. Using a relation between the hexagonal and the $\left(3.12^{2}\right)$ lattices, we also give bounds for the connective constant of the latter lattice.


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## 1. Introduction

In this work we give improved bounds for the connective constant of the hexagonal lattice. However, the methods used are valid for a larger class of lattices.

The main motivation is to improve the partial order induced by strict bounds for connective constants for different lattices, studied in [3]. In order to separate the hexagonal lattice from the $\left(4.8^{2}\right)$ lattice, we needed to apply a non-standard application of Kesten's method, which motivated a separate paper on the hexagonal lattice. A related partial order defined by percolation thresholds is studied in [10, 12].

### 1.1. Self-avoiding walks

A walk of length $n$ on a lattice is an alternating sequence of vertices and edges $\left\{v_{0}, e_{1}, v_{1}, e_{2}, \ldots, e_{n}, v_{n}\right\}$ such that the edge $e_{i}$ connects the vertices $v_{i-1}$ and $v_{i}$. The walk is self-avoiding if all vertices $v_{0}, v_{1}, \ldots, v_{n}$ are distinct.

For a regular graph, let $f_{n}$ denote the number of self-avoiding walks, starting at a fixed vertex. Hammersley [6] proved that there exists a constant $\mu$, called the connective constant of the lattice, such that

$$
\lim _{n \rightarrow \infty} f_{n}^{1 / n}=\mu
$$

Define the generating function for self-avoiding walks, sometimes called the susceptibility, by ( $f_{0}=1$ by convention)

$$
F(x)=\sum_{n=0}^{\infty} f_{n} x^{n}
$$

The generating function has radius of convergence $x_{c}=\mu^{-1}$.
The connective constant is unknown for all (truly) two-dimensional lattices, although Nienhuis [9] has presented strong evidence through non-rigorous computations that the connective constant for the hexagonal lattice equals $\sqrt{2+\sqrt{2}} \approx 1.847759$.

### 1.2. Bounds for the connective constant

Since the connective constants are unknown for most lattices, many papers concern bounds for the constants.

The first bounds for the connective constant, $\mu=\mu_{\text {hex }}$, of the hexagonal lattice were given by Fisher and Sykes [5] in 1959. They obtained the bounds

$$
1.7872<\mu<1.9276
$$

and also enumerated $f_{n}$, for $n \leqslant 20$. Sykes et al [11] extended the enumeration to $n \leqslant 34$ in 1972.

In his pioneering paper, Hammersley [6] showed that

$$
\mu<f_{n}^{1 / n}
$$

which, using $f_{34}$ gives

$$
\mu<f_{34}^{1 / 34}<1.9232
$$

although Sykes et al did not give this bound in their paper.
Ahlberg and Janson [1] used the above enumeration, and the fact that for the hexagonal lattice

$$
\begin{equation*}
\mu<\left(f_{n} / f_{2}\right)^{1 /(n-2)} \tag{1}
\end{equation*}
$$

to show that

$$
\mu<\left(f_{34} / f_{2}\right)^{1 / 32}<1.895 .
$$

In 1993, Alm [2] showed that

$$
\begin{equation*}
\mu<\left(\lambda_{1}(G(m, n))\right)^{1 /(n-m)} \tag{2}
\end{equation*}
$$

where $\lambda_{1}$ denotes the largest eigenvalue of the matrix $G(m, n)=\left(g_{i j}\right)_{K_{m} \times K_{m}}$, where the element $g_{i j}$ is the number of $n$-stepped self-avoiding walks that start with $\gamma_{i}$ and end with a translation of $\gamma_{j}$, and $\gamma_{1}, \ldots, \gamma_{K_{m}}$ are the different $m$-stepped self-avoiding walks (after taking symmetry considerations into account).

Using (2) with $n=34, m=12$, gave $K_{12}=736$, and the bound

$$
\mu<1.87603
$$

which is the currently best available upper bound for $\mu$.
To summarize, we know that

$$
1.7872<\mu<1.87603
$$

which should be compared with the supposedly correct value given by Nienhuis [9],

$$
\mu^{(N)}=\sqrt{2+\sqrt{2}} \approx 1.847759
$$

The bounds differ by $-0.061(-3.3 \%)$ and $+0.028(+1.5 \%)$ from this value, so that there is more room for improvement of the lower bound.

## 2. Upper bounds

Improved upper bounds are obtained by the method of Alm [2]. The improvement of computers in the last ten years makes it possible to handle much larger matrices as well as much longer self-avoiding walks.

Using (2) with $n=45, m=17$, we get $K_{17}=17700$, and the bound

$$
\mu<1.868832
$$

This reduces the difference between the upper bound and $\mu^{(N)}=\sqrt{2}+\sqrt{2}$ to $+0.021(+1.1 \%)$.
The computation took 944 CPU hours on a 1 GHz PC.
From the matrix $G(m, n)$, we can also compute $f_{n}$, thereby extending the enumeration of $f_{n}$ to $n \leqslant 45$. This was further extended to $n \leqslant 48$ by direct computations; $f_{48}$ requiring 928 CPU hours.

The values of $f_{n}, n \leqslant 48$ are given in table 1 together with upper bounds for $\mu$ for $m=17$ and $34 \leqslant n \leqslant 45$. For comparison, we also include the bounds (1) for $m=2$, i.e. $\bar{\mu}_{2}=\left(f_{n} / f_{2}\right)^{1 /(n-2)}$, which are the best bounds that can be obtained directly from the enumeration.

## 3. Lower bounds

### 3.1. The method of Kesten

In [8], Kesten presents a method for finding lower bounds for the connective constant, based on so called irreducible bridges.

Given a fixed embedding of the lattice in the plane, let the coordinates for a vertex $v$ be denoted by $(v(x), v(y))$. A bridge of length $n$ is a self-avoiding walk such that

$$
v_{0}(y)<v_{i}(y) \leqslant v_{n}(y) \quad \text { for } \quad i=1, \ldots, n-1
$$

Denote the number of bridges of length $n$ by $b_{n}$, and the generating function for bridges by ( $b_{0}=1$ )

$$
B(x)=\sum_{n=0}^{\infty} b_{n} x^{n}
$$

An irreducible bridge is a bridge that cannot be decomposed into two bridges. Denote the number of irreducible bridges of length $n$ by $a_{n}$, and the generating function for irreducible bridges by $\left(a_{0}=0\right)$

$$
A(x)=\sum_{n=1}^{\infty} a_{n} x^{n}
$$

Kesten proved that the connective constants for bridges and irreducible bridges are the same as for self-avoiding walks,

$$
\lim _{n \rightarrow \infty} b_{n}^{1 / n}=\lim _{n \rightarrow \infty} a_{n}^{1 / n}=\lim _{n \rightarrow \infty} f_{n}^{1 / n}=\mu
$$

Further, $A(x)$ and $B(x)$ are related by

$$
B(x)=\frac{1}{1-A(x)}
$$

Then, since $A(x)$ and $B(x)$ both have radius of convergence $\mu^{-1}$, it follows [8, theorem 5] that the solution of $A(x)=1$ is $\mu^{-1}$.

Table 1. Number of self-avoiding walks, $f_{n}$, and upper bounds, $\bar{\mu}_{m}$, for $m=2$ and $m=17$.

| $n$ | $f_{n}$ | $\bar{\mu}_{2}$ | $\bar{\mu}_{17}$ |
| :---: | :---: | :---: | :---: |
| 1 | 3 |  |  |
| 2 | 6 |  |  |
| 3 | 12 | 2.000000 |  |
| 4 | 24 | 2.000000 |  |
| 5 | 48 | 2.000000 |  |
| 6 | 90 | 1.967990 |  |
| 7 | 174 | 1.961010 |  |
| 8 | 336 | 1.955982 |  |
| 9 | 648 | 1.952042 |  |
| 10 | 1218 | 1.942840 |  |
| 11 | 2328 | 1.939314 |  |
| 12 | 4416 | 1.935031 |  |
| 13 | 8388 | 1.931770 |  |
| 14 | 15780 | 1.927509 |  |
| 15 | 29892 | 1.924934 |  |
| 16 | 56268 | 1.921862 |  |
| 17 | 106200 | 1.919545 |  |
| 18 | 199350 | 1.916866 |  |
| 19 | 375504 | 1.914895 |  |
| 20 | 704304 | 1.912692 |  |
| 21 | 1323996 | 1.910950 |  |
| 22 | 2479692 | 1.909029 |  |
| 23 | 4654464 | 1.907493 |  |
| 24 | 8710212 | 1.905836 |  |
| 25 | 16328220 | 1.904467 |  |
| 26 | 30526374 | 1.902999 |  |
| 27 | 57161568 | 1.901771 |  |
| 28 | 106794084 | 1.900472 |  |
| 29 | 199788408 | 1.899364 |  |
| 30 | 372996450 | 1.898197 |  |
| 31 | 697217994 | 1.897191 |  |
| 32 | 1300954248 | 1.896140 |  |
| 33 | 2430053136 | 1.895223 |  |
| 34 | 4531816950 | 1.894268 | 1.872434 |
| 35 | 8459583678 | 1.893427 | 1.872094 |
| 36 | 15769091448 | 1.892556 | 1.871650 |
| 37 | 29419727280 | 1.891782 | 1.871334 |
| 38 | 54816035922 | 1.890984 | 1.870930 |
| 39 | 102216080286 | 1.890269 | 1.870635 |
| 40 | 190380602052 | 1.889534 | 1.870266 |
| 41 | 354843312276 | 1.888871 | 1.869989 |
| 42 | 660671299170 | 1.888191 | 1.869650 |
| 43 | 1230891734724 | 1.887575 | 1.869390 |
| 44 | 2291023353264 | 1.886944 | 1.869077 |
| 45 | 4266787588320 | 1.886370 | 1.868832 |
| 46 | 7939282155480 | 1.885783 |  |
| 47 | 14780995214220 | 1.885246 |  |
| 48 | 27495750661500 | 1.884698 |  |



Figure 1. The hexagonal lattice, with an irreducible bridge of height 4.

Lower bounds for $\mu$ follow from the observation that, if $0 \leqslant \tilde{a}_{n} \leqslant a_{n}, n=2,3, \ldots, \infty$, then the reciprocal $x_{c}^{-1}$ of the solution $x_{c}$ to

$$
\sum_{n=1}^{\infty} \tilde{a}_{n} x^{n}=1
$$

is a lower bound for the connective constant $\mu$. Note that by letting $\tilde{a}_{n}=0$ for $n>N$, this allows us to truncate the infinite series.

Consider the standard embedding of the hexagonal lattice, as shown in figure 1. Define the height of a walk from $v_{0}$ to $v_{n}$ as the minimum number of vertical edges in any walk from $v_{0}$ to $v_{n}$. An irreducible bridge of length 22 and height 4 is shown in figure 1 . Let $b_{n}^{N}$ and $a_{n}^{N}$ denote the number of bridges and irreducible bridges of length $n$ and height $N$, and denote their generating functions by $B_{N}(x)$ and $A_{N}(x)$. Since $\sum_{k=1}^{\infty} a_{n}^{k}=a_{n}$, the reciprocal of the solution to

$$
\sum_{k=1}^{N} A_{k}(x)=\sum_{n=2}^{\infty}\left(\sum_{k=1}^{N} a_{n}^{k}\right) x^{n}=1
$$

is a lower bound for the connective constant $\mu$.
Note that irreducible bridges must end at a top vertex of a hexagon, and that an irreducible bridge of height $N \geqslant 2$ must use at least three vertical edges of each height except the first, and have length at least $6 N-2$. For height 1 , there are two irreducible bridges of each even length $n \geqslant 2$ (an irreducible bridge on the hexagonal lattice must have length at least 2).

Also, as every bridge either is an irreducible bridge, or can be decomposed into one irreducible bridge and one bridge, the following relation holds

$$
B_{N}(x)=A_{N}(x)+A_{N-1}(x) B_{1}(x)+A_{N-2}(x) B_{2}(x)+\cdots+A_{1}(x) B_{N-1}(x) .
$$

The generating function $A_{N}(x)$ can thus be obtained from $A_{k}(x), k=1, \ldots, N-1$ and $B_{k}(x), k=1, \ldots, N$, by
$A_{N}(x)=B_{N}(x)-A_{1}(x) B_{N-1}(x)-A_{2}(x) B_{N-2}(x)-\cdots-A_{N-1}(x) B_{1}(x)$.
In the next section we will see that, using methods from Alm and Janson [4], it is theoretically possible to calculate $B_{N}(x)$, and thus $A_{N}(x)$, in finite time for any fixed height $N$.

### 3.2. One-dimensional self-avoiding walks and bridges

In [4], Alm and Janson have studied self-avoiding walks on one-dimensional lattices. For our purpose, it is sufficient to define a one-dimensional lattices as a horizontal strip of finite height of a two-dimensional lattice.


Figure 2. The embedding of the hexagonal lattice used for finding the generating function for bridges of height 3 .

They show how to exactly compute the generating function for self-avoiding walks. The generating function is in principle obtained as one element in the inverse of a large matrix.

Only minor straightforward modifications are necessary to get the same result for the generating function for bridges.

To state the result precisely, we first need several definitions. Assume that the onedimensional lattice is embedded in the plane in such a way that it consists of an alternating sequence of isomorphic hinges and isomorphic sections.

A hinge consists of vertices, all with the same $x$-coordinate, and all edges between these vertices. A section consists of all edges between two adjacent hinges. For the hexagonal lattice, we use embeddings of the type showed in figure 2, in which the hinges consist of the vertices and the vertical edges. The sections consist of the dotted edges. Note that in the graph in figure 2 , self-avoiding walks that start at one of the bottom vertices, and ends at one of the top vertices, will be bridges of height 3 , according to the usual embedding of the hexagonal lattice.

Consider a bridge on the lattice. The appearance of the bridge in a given section is called a configuration. The appearance of the bridge in a given hinge is called a shape. The shape of a particular hinge is not necessarily determined by the adjacent configurations, but a shape completely determines the configurations on both sides.

Let $\Sigma$ be the set of all possible configurations, extended with two empty configurations, $\phi_{L}$ and $\phi_{R}$. Any finite walk must by convention start with the configuration $\phi_{L}$, and end with the configuration $\phi_{R}$. Any finite bridge must start at the lowest level, and end at the highest level.

There is a one-to-one correspondence between bridges and correctly connected alternating sequences $\left\{\phi_{L}=c_{0}, s_{1}, c_{1}, \ldots, s_{m}, c_{m}=\phi_{R}\right\}$ of configurations and shapes, correctly connected meaning that for every shape $s_{i}, c_{i-1}$ and $c_{i}$ are the configurations determined by $s_{i}$.

The generating function $B_{N}(x)$ for bridges of some fixed height $N$ is obtained from the following two square matrices, both indexed by $\Sigma . H=H(x)$ is a diagonal matrix with elements $x^{h_{i}}$, where $h_{i}$ is the number of edges in configuration $i . V=V(x)=\left\{v_{i j}\right\}$ has elements $v_{i j}=\sum x^{v_{k}}$, where the sum is over all shapes $k$ that can connect configuration $i$ to the left, with configuration $j$ to the right, and $v_{k}$ is the number of edges in shape $k$. By convention, $v_{i \phi_{L}}=v_{\phi_{R} i}=0$, for all $i$. Define the generating matrix $G=G(x)$ by $G=H V$.

Theorem 1 (Alm, Janson).

$$
B_{N}(x)=V_{\phi_{L}, \phi_{R}}+(V G)_{\phi_{L}, \phi_{R}}+\left(V G^{2}\right)_{\phi_{L}, \phi_{R}}+\cdots=\left(V(I-G)^{-1}\right)_{\phi_{L}, \phi_{R}} .
$$

Remark 1. A more straightforward way to get lower bounds from one-dimensional lattices would be to directly compute the connective constants for self-avoiding walks. However, this gives worse bounds than the method based on bridges, for one-dimensional lattices of the same heights.

### 3.3. Examples

We now illustrate theorem 1 with two examples, by calculating the generating function for irreducible bridges of heights 1 and 2.

For height 1 , there are four possible configurations, $\Sigma=\left\{\phi_{L}, \nwarrow, \searrow, \phi_{R}\right\}$, with lengths $0,1,1,0$. The generating matrix $G_{1}$ is given by
$G_{1}=H_{1} V_{1}=\left(\begin{array}{cccc}1 & 0 & 0 & 0 \\ 0 & x & 0 & 0 \\ 0 & 0 & x & 0 \\ 0 & 0 & 0 & 1\end{array}\right)\left(\begin{array}{cccc}0 & 1 & x^{2} & x^{2} \\ 0 & x & 0 & x \\ 0 & 0 & x & x \\ 0 & 0 & 0 & 0\end{array}\right)=\left(\begin{array}{cccc}0 & 1 & x^{2} & x^{2} \\ 0 & x^{2} & 0 & x^{2} \\ 0 & 0 & x^{2} & x^{2} \\ 0 & 0 & 0 & 0\end{array}\right)$.
The generating function is given by

$$
\begin{aligned}
A_{1}(x) & =B_{1}(x)=\left(V_{1}\left(I-G_{1}\right)^{-1}\right)_{1,4}=\frac{2 x^{2}}{1-x^{2}} \\
& =2 x^{2}+2 x^{4}+2 x^{6}+2 x^{8}+\cdots .
\end{aligned}
$$

For height 2, there are eight configurations,

$$
\Sigma=\left\{\phi_{L}, \stackrel{\searrow}{\nwarrow}, \nwarrow, \nwarrow, \downarrow, \searrow, \phi_{R}, \quad \searrow\right\}
$$

The matrices $H_{2}$ and $V_{2}$ are given by
$H_{2}=\left(\begin{array}{cccccccc}1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & x^{2} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & x & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & x & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & x & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & x & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & x^{2}\end{array}\right) \quad V_{2}=\left(\begin{array}{cccccccc}0 & x^{2} & 1 & x^{2} & x^{4} & x^{2} & x^{4} & x^{2} \\ 0 & x^{2} & 0 & x^{2} & x^{2} & x^{2} & x^{2} & x^{2} \\ 0 & 0 & x & x & 0 & 0 & x^{3} & x^{3} \\ 0 & 0 & x^{3} & x & 0 & 0 & x & 0 \\ 0 & 0 & 0 & 0 & x & x & x & x \\ 0 & 0 & 0 & 0 & x^{3} & x & x^{3} & x \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & x^{2} & x^{2}\end{array}\right)$.
The generating function is

$$
\begin{aligned}
A_{2}(x) & =B_{2}(x)-A_{1}(x) B_{1}(x)=\left(V_{2}\left(I-G_{2}\right)^{-1}\right)_{1,7}-A_{1}(x)^{2} \\
& =\frac{2 x^{10}\left(x^{4}+4 x^{2}+4\right)}{\left(1-x^{2}\right)^{2}\left(1+x^{2}\right)\left(1-x^{2}-x^{4}-x^{8}\right)}=8 x^{10}+24 x^{12}+58 x^{14}+\cdots
\end{aligned}
$$

### 3.4. Summary

In this section we summarize the method used for the lower bound. Assume that we have calculated the generating functions $B_{k}$ for bridges of height $k=1, \ldots, N$. From these we can calculate $A_{k}$, by equation (3). A lower bound for the connective constant $\mu$ is achieved by the reciprocal $x_{c}^{-1}$ of the solution $x_{c}$ to the equation $\sum_{k=1}^{N} A_{k}(x)=1$.

Note that, if we underestimate the functions $B_{k}$, we get underestimates of the functions $A_{k}$, and therefore, still get a lower bound. It is therefore permissible to truncate the infinite series $B_{k}$, for the cases where it is feasible to carry out the symbolic inversion of the matrix, and also to truncate the infinite series representation of the inverse matrix.

Table 2. Number of irreducible bridges.

| Height |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $n$ | 1 | 2 | 3 | 4 | 5 | 6 | 7 |
| 2 | 2 |  |  |  |  |  |  |
| 4 | 2 |  |  |  |  |  |  |
| 6 | 2 |  |  |  |  |  |  |
| 8 | 2 |  |  |  |  |  |  |
| 10 | 2 | 8 |  |  |  |  |  |
| 12 | 2 | 24 |  |  |  |  |  |
| 14 | 2 | 58 |  |  |  |  |  |
| 16 | 2 | 116 | 40 |  |  |  |  |
| 18 | 2 | 226 | 248 |  |  |  |  |
| 20 | 2 | 418 | 956 |  |  |  |  |
| 22 | 2 | 764 | 2932 | 232 |  |  |  |
| 24 | 2 | 1368 | 8158 | 2208 |  |  |  |
| 26 | 2 | 2438 | 21194 | 11908 |  |  |  |
| 28 | 2 | 4312 | 52768 | 48924 | 1456 |  |  |
| 30 | 2 | 7612 | 127424 | 174384 | 18656 |  |  |
| 32 | 2 | 13398 | 301336 | 567066 | 130180 |  |  |
| 34 | 2 | 23564 | 701240 | 1734242 | 669868 | 9584 |  |
| 36 | 2 | 41398 | 1613176 | 5077228 | 2896786 | 154112 |  |
| 38 | 2 | 72708 | 3678486 | 14399854 | 11186920 | 1321320 |  |
| 40 | 2 | 127646 | 8332878 | 39863556 | 39960038 | 8176732 | 65136 |
| 42 | 2 | 224070 | 18781132 | 108330740 | 134910038 | 41577438 | 1258416 |
| 44 | 2 | 393274 | 42167658 | 290120946 | 436772298 | 185506724 | 12797828 |
| 46 | 2 | 690222 | 94393440 | 768000442 | 1369384066 | 754670432 | 92575300 |
| 48 | 2 | 1211320 | 210817326 | 2014040116 | 4287469968 | 2868092528 | 541473298 |
| 50 | 2 | 2125800 | 469995164 | 5241551508 | 12553608776 | 10351754668 | 2740477924 |
| 52 | 2 | 3730590 | 1046346578 | 13555798878 | 37039148380 | 35894677308 | 12497716690 |
| 54 | 2 | 6546818 | 2326934596 | 34876762888 | 107870357502 | 120587726960 | 52714512594 |
| 56 | 2 | 11488942 | 5170378690 | 89344535622 | 310799090686 | 394963366374 | 209362196128 |
| 58 | 2 | 20161784 | 11480731734 | 228047517858 | 887494845612 | 1267269921544 | 793008736988 |
| 60 | 2 | 35381548 | 25479343670 | 580303521948 | 2515218274794 | 3998003947894 | 2891762035512 |
| 62 | 2 | 62090392 | 56523233522 | 1472856568902 | 7082773418692 | 12437565678140 | 10224629512728 |
| 64 | 2 | 108961132 | 125349924212 | 3729993423200 | 19835888844516 | 38241846805060 | 35247308133632 |
| 66 | 2 | 191213568 | 277913356354 | 9428372543664 | 55290614198678 | 116425397228820 | 118980531760554 |
| 68 | 2 | 335556516 | 616036743246 | 23793707134584 | 153488979691246 | 351481068641732 | 394632167359492 |
| 70 | 2 | 588860754 | 1365319252836 | 59962719844118 | 424578943200246 | 1053470052659232 | 1289668249938810 |
| 72 | 2 | 1033378688 | 3025567446724 | 150930254650748 | 1170815299517686 | 3137874062518550 | 4162060662237216 |
| 74 | 2 | 1813453306 | 6704031069242 | 379504185427756 | 3219812806690196 | 9295979825494524 | 13288671828545886 |
| 76 | 2 | 3182388814 | 14853570233752 | 953369709263106 | 8833331985285592 | 27409160402323884 | 42039183397649906 |
| 78 | 2 | 5584703188 | 32907775865780 | 2393102745030608 | 24181849819174446 | 80479402119427078 | 131938075461499130 |
| 80 | 2 | 9800471012 | 72902886657418 | 6002886315196280 | 66073638249791486 | 235435386138206528 | 411227122483745614 |
| 82 | 2 | 17198627838 | 161500519239580 | 15048570170389500 | 180230973427403650 | 686491758206462152 | 1273995129812876972 |
| 84 | 2 | 30181488004 | 357758261638838 | 37705022115819458 | 490874221247694464 | 1995853645048476200 | 3925967882491042342 |

### 3.5. Results

For heights up to 4 , we were able to find exact expressions for the generating function for irreducible bridges. These are given in section 3.6. For heights 5, 6 and 7, we were also able to determine the generating matrix $G$, but were unable to carry out the matrix inversion. Instead we truncated the infinite series representation of the inverse, by an iterative method described in section 3.7.

The number of irreducible bridges for heights up to 7 , and lengths up to 84 are given in table 2.

To improve the lower bound, we have also computed the number of irreducible bridges of heights 8,9 and 10 , for lengths up to 58,

$$
\begin{aligned}
& \begin{array}{l}
A_{8}(x)=452960 x^{46}+10207872 x^{48}+120066252 x^{50}+99411028 x^{52}+657310510 x^{54} \\
\quad+37244196500 x^{56}+188193918866 x^{58}+O\left(x^{60}\right) \\
A_{9}(x)=3204784 x^{52}+82465376 x^{54}+1100970508 x^{56}+10265702604 x^{58}+O\left(x^{60}\right) \\
A_{10}(x)=22982032 x^{58}+O\left(x^{60}\right) .
\end{array} .
\end{aligned}
$$

The best lower bound found in this work is 1.833009 764. The difference between this bound and $\mu^{(N)}=\sqrt{2+\sqrt{2}}$ is $-0.015(-0.8 \%)$, slightly smaller than the difference for the upper bound.

Remark 2. The computations for $n=58$ were distributed over a large number of computers. The total computation time was 22400 CPU hours ( 1.6 GHz ).

### 3.6. The generating functions

Below, we give the generating function for heights up to 4. $D_{N}$ and $N_{N}$ denotes the denominator and numerator of $A_{N}$.

$$
A_{1}(x)=\frac{2 x^{2}}{1-x^{2}}
$$

$$
A_{2}(x)=\frac{2 x^{10}\left(x^{4}+4 x^{2}+4\right)}{\left(1-x^{2}\right)^{2}\left(1+x^{2}\right)\left(1-x^{2}-x^{4}-x^{8}\right)}
$$

$$
\begin{aligned}
& D_{3}(x)=\left(x^{4}-1\right)^{2}\left(x^{8}-2 x^{6}+3 x^{4}-3 x^{2}+1\right)\left(x^{38}-5 x^{36}+14 x^{34}-27 x^{32}+36 x^{30}-34 x^{28}\right. \\
&+19 x^{26}-2 x^{24}-6 x^{22}+11 x^{20}-18 x^{18}+27 x^{16}-23 x^{14} \\
&\left.+9 x^{12}-x^{10}+4 x^{6}-8 x^{4}+5 x^{2}-1\right) \\
& N_{3}(x)=-2\left(x^{38}+4 x^{36}-20 x^{34}+45 x^{32}-57 x^{30}-6 x^{28}+127 x^{26}-243 x^{24}\right. \\
&+205 x^{22}-22 x^{20}-54 x^{18}-21 x^{16}+137 x^{14}-2 x^{12} \\
&\left.-209 x^{10}+147 x^{8}+38 x^{6}-34 x^{4}-36 x^{2}+20\right) x^{16}
\end{aligned}
$$

$$
\begin{aligned}
D_{4}(x)=\left(x^{2}+\right. & +1)^{9}\left(x^{5}-x^{4}+x^{2}-x+1\right)^{2}\left(x^{5}+x^{4}-x^{2}-x-1\right)^{2}\left(x^{4}+1\right)^{2}\left(x^{40}-4 x^{38}\right. \\
& +4 x^{36}+5 x^{34}-17 x^{32}+23 x^{30}-16 x^{28}-4 x^{26}+17 x^{24}-15 x^{22}+8 x^{20} \\
& \left.+5 x^{18}-14 x^{16}+5 x^{14}+2 x^{12}-2 x^{10}+4 x^{8}-x^{6}-4 x^{4}+1\right)^{2}\left(x^{37}-x^{36}\right. \\
& -4 x^{35}+6 x^{34}+6 x^{33}-15 x^{32}-4 x^{31}+19 x^{30}+4 x^{29}-14 x^{28}-12 x^{27} \\
& +13 x^{26}+17 x^{25}-19 x^{24}-13 x^{23}+21 x^{22}+14 x^{21}-19 x^{20}-18 x^{19}+17 x^{18} \\
& +10 x^{17}-14 x^{16}-6 x^{15}+15 x^{14}+8 x^{13}-10 x^{12}-5 x^{11}+4 x^{10}+4 x^{9}-8 x^{8} \\
& \left.-3 x^{7}+5 x^{6}+x^{5}-x^{4}-x^{3}+2 x^{2}-1\right)^{2}\left(x^{37}+x^{36}-4 x^{35}-6 x^{34}+6 x^{33}\right. \\
& +15 x^{32}-4 x^{31}-19 x^{30}+4 x^{29}+14 x^{28}-12 x^{27}-13 x^{26}+17 x^{25}+19 x^{24} \\
& -13 x^{23}-21 x^{22}+14 x^{21}+19 x^{20}-18 x^{19}-17 x^{18}+10 x^{17}+14 x^{16}-6 x^{15} \\
& -15 x^{14}+8 x^{13}+10 x^{12}-5 x^{11}-4 x^{10}+4 x^{9}+8 x^{8}-3 x^{7}-5 x^{6}+x^{5} \\
& \left.+x^{4}-x^{3}-2 x^{2}+1\right)^{2}(x-1)^{12}(x+1)^{12}\left(x^{8}-x^{6}+x^{4}+x^{2}-1\right)^{2}\left(x^{8}-x^{6}\right. \\
& \left.-x^{4}+x^{2}+1\right)^{2}\left(x^{11}+x^{10}-x^{9}-2 x^{8}+x^{6}-x^{5}+x^{3}+x^{2}-x-1\right)^{2} \\
& \times\left(x^{11}-x^{10}-x^{9}+2 x^{8}-x^{6}-x^{5}+x^{3}-x^{2}-x+1\right)^{2} \\
& \times\left(x^{3}+x^{2}-1\right)^{4}\left(x^{3}-x^{2}+1\right)^{4}
\end{aligned}
$$

$$
\begin{aligned}
& N_{4}(x)=-2\left(x^{11}+x^{10}-x^{9}-2 x^{8}+x^{6}-x^{5}+x^{3}+x^{2}-x-1\right)\left(x^{11}-x^{10}-x^{9}+2 x^{8}\right. \\
&\left.-x^{6}-x^{5}+x^{3}-x^{2}-x+1\right)\left(x^{37}+x^{36}-4 x^{35}-6 x^{34}+6 x^{33}+15 x^{32}\right. \\
&-4 x^{31}-19 x^{30}+4 x^{29}+14 x^{28}-12 x^{27}-13 x^{26}+17 x^{25}+19 x^{24}-13 x^{23} \\
&-21 x^{22}+14 x^{21}+19 x^{20}-18 x^{19}-17 x^{18}+10 x^{17}+14 x^{16}-6 x^{15}-15 x^{14}
\end{aligned}
$$

$$
\begin{aligned}
& +8 x^{13}+10 x^{12}-5 x^{11}-4 x^{10}+4 x^{9}+8 x^{8}-3 x^{7}-5 x^{6}+x^{5}+x^{4}-x^{3} \\
& \left.-2 x^{2}+1\right)\left(x^{37}-x^{36}-4 x^{35}+6 x^{34}+6 x^{33}-15 x^{32}-4 x^{31}+19 x^{30}+4 x^{29}\right. \\
& -14 x^{28}-12 x^{27}+13 x^{26}+17 x^{25}-19 x^{24}-13 x^{23}+21 x^{22}+14 x^{21}-19 x^{20} \\
& -18 x^{19}+17 x^{18}+10 x^{17}-14 x^{16}-6 x^{15}+15 x^{14}+8 x^{13}-10 x^{12}-5 x^{11} \\
& \left.+4 x^{10}+4 x^{9}-8 x^{8}-3 x^{7}+5 x^{6}+x^{5}-x^{4}-x^{3}+2 x^{2}-1\right)\left(x^{198}-5 x^{196}\right. \\
& -94 x^{194}+1439 x^{192}-10077 x^{190}+45706 x^{188}-146052 x^{186}+322139 x^{184} \\
& -375099 x^{182}-460084 x^{180}+3735706 x^{178}-10989654 x^{176} \\
& +21041143 x^{174}-26443197 x^{172}+13762470 x^{170}+26655182 x^{168} \\
& -83450197 x^{166}+116277400 x^{164}-79099657 x^{162}-26911408 x^{160} \\
& +117139658 x^{158}-66719812 x^{156}-163339763 x^{154}+420127625 x^{152} \\
& -414373898 x^{150}-40794436 x^{148}+792055755 x^{146}-1369812592 x^{144} \\
& +1339768016 x^{142}-713686008 x^{140}+13379549 x^{138}+149203153 x^{136} \\
& +366222831 x^{134}-1070571661 x^{132}+1286138525 x^{130}-776725830 x^{128} \\
& -56615258 x^{126}+595347458 x^{124}-586136637 x^{122}+301833490 x^{120} \\
& -174524358 x^{118}+329622043 x^{116}-497688054 x^{114}+351910601 x^{112} \\
& +110012919 x^{110}-541133599 x^{108}+598537662 x^{106}-284947505 x^{104} \\
& -81751322 x^{102}+193056310 x^{100}-23672554 x^{98}-188554083 x^{96} \\
& +218784743 x^{94}-61217653 x^{92}-131592311 x^{90}+207344673 x^{88} \\
& -139738148 x^{86}+41821218 x^{84}-532196 x^{82}+1689902 x^{80} \\
& -25387860 x^{78}+36161567 x^{76}+1501343 x^{74}-30406174 x^{72} \\
& +12099108 x^{70}+5503781 x^{68}-11955348 x^{66}+12806700 x^{64} \\
& +5866625 x^{62}-13834418 x^{60}+1869202 x^{58}+1479416 x^{56} \\
& -6953954 x^{54}+12983920 x^{52}-1064824 x^{50}-5192613 x^{48} \\
& +61083 x^{46}-2124568 x^{44}+471817 x^{42}+4762797 x^{40}-9273 x^{38} \\
& -2486033 x^{36}-1428960 x^{34}+153688 x^{32}+1762191 x^{30}+549830 x^{28} \\
& -877409 x^{26}-560376 x^{24}+220009 x^{22}+275271 x^{20}+54321 x^{18} \\
& -115656 x^{16}-49531 x^{14}+30788 x^{12}+19929 x^{10}-7606 x^{8}-3448 x^{6} \\
& \left.+974 x^{4}+404 x^{2}-116\right)\left(x^{3}-x^{2}+1\right)^{2}\left(x^{3}+x^{2}-1\right)^{2} \\
& \times\left(x^{5}-x^{4}+x^{2}-x+1\right)^{2}\left(x^{5}+x^{4}-x^{2}-x-1\right)^{2}\left(x^{4}+1\right)^{2} \\
& \times\left(x^{8}-x^{6}-x^{4}+x^{2}+1\right)^{2}\left(x^{2}+1\right)^{5}(x-1)^{8}(x+1)^{8} x^{22}
\end{aligned}
$$

### 3.7. Algorithmic issues

We have calculated the generating matrix for bridges of heights up to 7. The computations were done by a computer program written in C++. The program iterates through all possible pairs of configurations, finds all possible shapes that connects the two configurations and generates the matrix in Maple format.

The program is quite time consuming, the case of height 7 required a total computation time of several weeks of CPU time on a standard desktop computer. However, the matrix was divided into submatrices, which were generated on separate computers.


Figure 3. The (3.12 ${ }^{2}$ ) lattice.

Table 3. Summary of data for the computations of the generating functions.

| Height | 1 | 2 | 3 | 4 | 5 | 6 | 7 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | ---: |
| Dimension of generating matrix | 4 | 8 | 18 | 44 | 118 | 338 | 1024 |
| Iteration steps | $\infty$ | $\infty$ | $\infty$ | $\infty$ | 130 | 100 | 42 |
| Exact enumeration up to term | $\infty$ | $\infty$ | $\infty$ | $\infty$ | 260 | 200 | 84 |

The matrices and generating functions were treated in Maple. For heights up to 4 Maple was able to handle the symbolic inversions. For heights 5,6 and 7 we calculated a truncation of the infinite series representation, by the following iterative method.

Let $J_{\phi_{R}}$ be the column vector with a one in position $\phi_{R}$, and zeros elsewhere, and let $G_{\phi_{R}}$ be column $\phi_{R}$ in $G$. Let $F^{(1)}=G_{\phi_{R}}+J_{\phi_{R}}$, and iteratively, for $k \geqslant 2, F^{(k)}=G F^{(k-1)}+J_{\phi_{R}}$. Then

$$
\left(V F^{(n)}\right)_{\phi_{L}}=V_{\phi_{L}, \phi_{R}}+(V G)_{\phi_{L}, \phi_{R}}+\left(V G^{2}\right)_{\phi_{L}, \phi_{R}}+\cdots+\left(V G^{n}\right)_{\phi_{L}, \phi_{R}} .
$$

For heights 5, 6 and 7, we performed 130, 100 and 42 iteration steps. For height 7, we are currently unable to do further iterations, due to local computer limitations. In table 3, the dimensions of the generating matrices, and the number of iteration steps are summarized.

It is easy to check that $n$ iteration steps give the correct number of bridges, and thus also irreducible bridges, for lengths up to $2 n$. The lower bound is further improved by including all positive higher order terms.

## 4. Bounds for the $\left(3.12^{2}\right)$ lattice

Jensen and Guttmann [7] show, by studying the generating functions for self-avoiding polygons, that the connective constant of the $\left(3.12^{2}\right)$ lattice, see figure 3 , is related to the connective constant of the hexagonal lattice through

$$
\begin{equation*}
\frac{1}{\mu_{\mathrm{hex}}}=\frac{1}{\mu_{\left(3.12^{2}\right)}^{2}}+\frac{1}{\mu_{\left(3.12^{2}\right)}^{3}} \tag{4}
\end{equation*}
$$

Using Nienhuis' value [9], $\mu^{(N)}=\sqrt{2+\sqrt{2}}$ for the hexagonal lattice gives

$$
\begin{aligned}
\mu_{\left(3.12^{2}\right)}^{(N)}= & \frac{12^{\frac{1}{3}} \cdot(2+\sqrt{2})^{\frac{1}{6}} \cdot(9+\sqrt{81-12 \sqrt{2+\sqrt{2}}})^{\frac{1}{3}}}{6} \\
& +\frac{12^{\frac{1}{3}} \cdot(2+\sqrt{2})^{\frac{1}{6}} \cdot(9-\sqrt{81-12 \sqrt{2+\sqrt{2}}})^{\frac{1}{3}}}{6} \approx 1.711041 .
\end{aligned}
$$

Using the relation (4) and the bounds for $\mu_{\text {hex }}$ from the previous sections, we get corresponding bounds for $\mu_{\left(3.12^{2}\right)}$

$$
1.705263<\mu_{\left(3.12^{2}\right)}<1.719254
$$

Here, the bounds differ from the presumably correct value $\mu_{\left(3.12^{2}\right)}^{(N)}$ by $-0.0058(-0.34 \%)$ and +0.0082 ( $+0.48 \%$ ), respectively.

Remark 3. The upper bound for $\mu_{\left(3.12^{2}\right)}$ obtained above is much better than the corresponding upper bound obtained by direct computations on the (3.12 ${ }^{2}$ ) lattice. Using (2) with $n=48$ and $m=18\left(K_{18}=23976\right)$ gave, after 620 CPU hours, the upper bound

$$
\mu_{\left(3.12^{2}\right)}<1.72922
$$

overestimating $\mu_{\left(3.12^{2}\right)}^{(N)}$ by 0.018 (1.1\%).
Remark 4. Jensen and Guttmann [7] show that the generating function for self-avoiding polygons on the $\left(3.12^{2}\right)$ lattice can be obtained from the corresponding generating function for the hexagonal lattice through the transformation $x \rightarrow x^{2}+x^{3}$ (apart from an initial term $x^{3} / 3$ corresponding to the triangles of $\left(3.12^{2}\right)$ ). This suffices to show the relation (4), as self-avoiding polygons have the same connective constant as self-avoiding walks.

However, Jensen and Guttmann also claim a similar relation between the generating functions for self-avoiding walks on these lattices, but their relation is incorrect, and it is doubtful whether such a relation exists.

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